

Aligned magnetohydrodynamic effect on magnetic nanoparticle with different base fluids past a moving inclined plate



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ABSTRACT

This paper deals with the numerical solutions for the aligned MHD free convection laminar boundary layer flow over a moving inclined plate for two magnetic nanofluids, namely Fe_3O_4 -water and Fe_3O_4 -kerosene. It is assumed that the left surface of the plate is in contact with a hot fluid while the cold fluid is on the right surface. The mathematical model has been constructed and based on the Tiwari-Das model, appropriate similarity transformations are used to convert the governing partial differential equations into nonlinear ordinary differential equations and solved numerically using the Keller-Box method. Numerical results for the skin friction coefficient and local Nusselt number were presented whilst the velocity and temperature profiles were illustrated graphically and analyzed. It is found that the velocity increases and temperature decrease with an increase of aligned magnetic field angle parameter, magnetic strength parameter, and Grashof number while the velocity decreases and temperature increase when inclined plate angle parameter and volume fractions of nanoparticles increase. For the convective parameter, both velocity and temperature profile increase when the Biot number increase. Comparisons with previously published studies are performed and excellent agreement is obtained.

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1. Introduction

Nanofluids have wide applications in various fields and have provided significant importance towards the enhancement of heat transfer. Their applications include desalination, cavity problem, solar thermal collector, nuclear system cooling, electronics cooling, solar water heating, heat exchanger, biomedicine, fuel cells, transportation (engine cooling/vehicle thermal management), and so on (Wu and Zhao, 2013; Hussein et al., 2014; Mitra, 2018). The term 'nanofluids' refers to a new class of engineered heat transfer fluids, which contain metallic particles with average particle sizes of 10nm and can be produced by current nanophase technology. The term was initially introduced by Choi and Eastman (1995).

Magnetic nanofluids consist of a colloidal mixture of magnetic nanoparticles with sizes in the range of 2–10nm and base liquid (Philip et al., 2007). The magnetic features of magnetic nanofluids are comparable to those of bulk magnetic materials and allow the retaining of common Newtonian fluids' flow characteristics. As indicated by several studies, a variety of techniques is applied to enhance heat transfer processes, especially for several ordinary fluids including water, toluene, ethylene glycol, and mineral oils. An innovative way to improve the heat transfer performance of fluids is to suspend nanophase particles in heating or cooling fluids (Xuan and Li, 2000; Alsaedi et al., 2012). Kamal et al. (2019) concentrated on the conjugate study of flow and heat transfer considering two physical effects such as g-jitter and thermal radiation on a three-dimensional stagnation point region. Recently, Nayan et al. (2022) used hybrid nanofluids in their study to increase heat transfer performance.

Magnetohydrodynamic (MHD) deals with the study of electrically conducting fluids that can move around in a magnetic field. An electric current is induced in the fluid when there is a change in the

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magnetic field that cuts the moving fluid. Rana et al. (2012) investigated an incompressible nanofluid's steady mixed convection boundary layer flow along an inclined plate that was embedded in a porous medium. In their study, it was observed that the Nusselt number decreased with a rise in the thermophoresis number or Brownian motion number, whereas increasing the plate's angle led to an increase in the Nusselt number. Research by Ilias et al. (2017) focused on the free convection boundary layer flow over an inclined plate for two ferrofluids, Fe_3O_4 -water and Fe_3O_4 -kerosene. The numerical solutions are carried out by the Keller-Box method. It can be concluded that the skin friction and heat transfer rate of both ferrofluids are increasing in all parameters except for the angle of the inclined plate. Anjali Devi and Suriyakumar (2013) examined the numerical solution of MHD mixed convection nanofluid flow over an inclined stretching plate by considering the heat generation and suction effects. Mohamad et al. (2017) conducted a study to investigate the unsteady MHD boundary layer flow of a Casson fluid in a porous medium past an arbitrary wall shear stress vertical plate with mass and heat transfer. Meanwhile, Noranuar et al. (2021a; 2021b) used two types of carbon nanotubes which are single-walled carbon nanotubes and multi-walled carbon nanotubes in non-coaxial rotation flow affected by magnetohydrodynamics and porosity to analyze the fluid flow and heat transfer. In recent years, numerous authors have investigated the issues associated with the MHD boundary layer flows and heat transfer by considering the different problems and situations (Sheikholeslami and Ganji, 2014; Reddy et al., 2017; Ullah et al., 2017; Ilias et al., 2018; 2020; Ahmad et al., 2019; Ismail et al., 2021; Bosli et al., 2022).

Many industrial applications, natural processes, and chemical processing systems are faced with free convection processes that involve a combination mechanism of mass and heat transfer. The convective boundary condition is mostly used to define a linear convective heat exchange condition for one or more algebraic entities in thermal. Ostrach (1952) who started the convection flow study, used an integral method to make a technical note on transient free convection flow's similarity solution past a semi-infinite vertical plate. Rawi et al. (2016; 2018) studied the effect of g-jitter on the mixed convection flow of Jeffrey fluid and micropolar nanofluids past an inclined stretching sheet with the presence of nanoparticles, while Hamdan et al. (2020) and Rosaidi et al. (2022) conducted their research based on free convection flow fluid since it is simpler to analyse the behaviors of the magnetic nanofluids because it naturally happens without affected by any external forces. Heat transfer analysis with convective boundary conditions is evoked in processes including high temperatures such as thermal energy storage, gas turbines, and nuclear plants (Akbar et al., 2013; Nadeem et al., 2013). There are several studies focused on the

effect of a convective boundary condition on boundary layer flow and heat transfer over a stretching or shrinking surface of nanofluid (Alsaedi et al., 2012; Makinde and Aziz, 2011; Yao et al., 2011; Ghazali et al., 2021). Recently, Kanafiah et al. (2022) focused on the problem of steady combined convective transport by considering an alternative development of heat transfer fluids to solve the performance of the heat transfer process in industrial sectors.

However, to the extent of the authors' knowledge, attempts are yet to be made to address the issues of aligned MHD free convection heat transfer flow regarding magnetic nanofluids over a moving inclined plate with convective boundary condition. Hence, a study was conducted in an attempt to understand free convection boundary layer flow over a moving inclined plate for two magnetic nanofluids, namely Fe_3O_4 -water and Fe_3O_4 -kerosene.

2. Mathematical formulation

Consider the hydromagnetic free convection flow of magnetic nanofluid to be steady, incompressible, two-dimensional, and laminar over a moving inclined plate with the aligned and transverse magnetic field. The semi-infinite plate is inclined at an acute angle γ to the vertical axis measured in the clockwise direction and situated in an otherwise quiescent ambient fluid at the temperature T_∞ . U_∞ is the constant free stream velocity? The plate is moving with constant velocity $U_w(x) = \lambda U_\infty$ where $U_w(x)$ the plate velocity is, λ is the plate velocity parameter. The physical coordinates (x, y) are chosen such that x -axis measured along the inclined plate and y -axis is normal to the surface of the plate. The plate is heated by convection through a hot fluid at the temperature T_f with heat transfer coefficient h_f . The gravitational acceleration \mathbf{g} is acting downward. The physical configuration of the problem is shown in Fig. 1. Water and kerosene are used as the base fluids with magnetite (Fe_3O_4) as nanoparticles. The base fluids and nanoparticles are in thermal equilibrium and no slip occurs between them. The spherical shaped nanoparticles are considered. The viscous dissipation and radiation are neglected in the analysis.

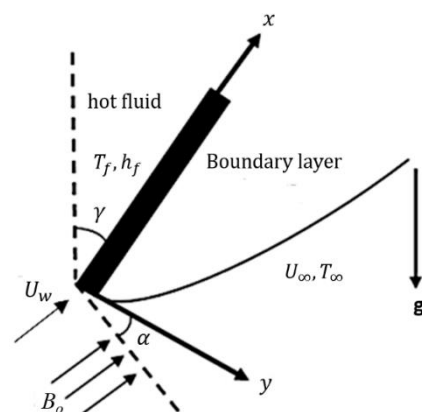


Fig. 1: Physical model and coordinate system

Under the above assumptions and following [Tiwari and Das \(2007\)](#), the equations of MHD boundary layer flow are:

$$\frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} = 0 \tag{1}$$

$$u \frac{\partial u}{\partial x} + v \frac{\partial u}{\partial y} = \frac{\mu_{nf}}{\rho_{nf}} \frac{\partial^2 u}{\partial y^2} + \frac{(\rho\beta)_{nf}}{\rho_{nf}} g \cos\gamma (T - T_\infty) - \frac{\sigma B^2(x)}{\rho_{nf}} \sin^2\alpha (u - U_\infty) \tag{2}$$

$$u \frac{\partial T}{\partial x} + v \frac{\partial T}{\partial y} = \alpha_{nf} \frac{\partial^2 T}{\partial y^2} \tag{3}$$

The boundary conditions for this problem are given by,

$$u = U_w(x) = \lambda U_\infty, \quad v = 0, \quad -k_{nf} \frac{\partial T}{\partial y} = h_f (T_f - T) \tag{4}$$

on

$$y = 0 \\ u \rightarrow U_\infty, \quad T \rightarrow T_\infty, \text{ as } y \rightarrow \infty$$

where, u and v are the x (along the inclined plate) and y (normal to the plate) components of velocities, respectively. T is the temperature of the magnetic nanofluid, σ is the electrical conductivity and α is the aligned magnetic field angle. The transverse magnetic field assumed to be a function of the distance from the origin is defined as $B(x) = B_0 x^{-\frac{1}{2}}$ with $B_0 \neq 0$, where x is the coordinate along the plate and B_0 is the magnetic field strength. The effective properties of magnetic nanofluids may be expressed in terms of the properties of base fluids, nanoparticles, and the volume fraction of solid nanoparticles as follow ([Khan et al., 2015](#)).

$$\rho_{nf} = (1 - \phi)\rho_f + \phi\rho_s, \quad \mu_{nf} = \frac{\mu_f}{(1 - \phi)^{2.5}}, \\ (\rho C_p)_{nf} = (1 - \phi)(\rho C_p)_f + \phi(\rho C_p)_s, \quad (\rho\beta)_{nf} = (1 - \phi)(\rho\beta)_f + \phi(\rho\beta)_s \tag{5}$$

$$\alpha_{nf} = \frac{k_{nf}}{(\rho C_p)_{nf}}, \quad \frac{k_{nf}}{k_f} = \frac{k_s + 2k_f - 2\phi(k_f - k_s)}{k_s + 2k_f + \phi(k_f - k_s)}$$

where, ρ_{nf} is the effective density, ϕ is the solid volume fraction, ρ_f is the density of the pure base fluid, ρ_s is the density of nanoparticles, μ_f is the dynamic viscosity of the base fluids, μ_{nf} is the effective dynamic viscosity, $(\rho C_p)_{nf}$ is the heat capacity of the magnetic nanofluids, $(\rho C_p)_f$ is the specific heat parameters of the base fluids, $(\rho C_p)_s$ is the specific heat parameters of nanoparticles, $(\rho\beta)_{nf}$ is the thermal expansion coefficient, α_{nf} is the thermal diffusivity of the magnetic nanofluids, k_{nf} is the thermal conductivity of the magnetic base fluid, k_f is the thermal conductivity of the magnetic nanofluids and k_s is the thermal conductivity of nanoparticles. Continuity 1 is satisfied by introducing a stream function $\psi(x, y)$ below

$$u = \frac{\partial \psi}{\partial y}, \quad v = -\frac{\partial \psi}{\partial x} \tag{6}$$

The following similarity variables are introduced:

$$\eta = y \sqrt{\frac{U_\infty}{\nu_f x}} = \frac{y}{x} \sqrt{Re_x}, \quad \psi = \nu_f \sqrt{Re_x} f(\eta), \quad \theta(\eta) = \frac{T - T_\infty}{T_f - T_\infty} \tag{7}$$

where, η is the similarity variable, $Re_x = U_\infty x / \nu_f$ is the Reynolds number, $f(\eta)$ the non-dimensional stream function and $\theta(\eta)$ the non-dimensional temperature.

On the use of Eqs. 5-7, Eqs. 2-3 reduce to the following nonlinear system of ordinary differentials equations:

$$f''' + (1 - \phi)^{2.5} \left(1 - \phi + \phi \left(\frac{\rho_s}{\rho_f} \right) \right) \frac{1}{2} f f'' \\ + (1 - \phi)^{2.5} M (1 - f') \sin^2\alpha + (1 - \phi)^{2.5} \left(1 - \phi + \phi \left(\frac{\rho\beta_s}{\rho\beta_f} \right) \right) Gr_x \theta \cos\gamma = 0 \tag{8}$$

$$\left(\frac{k_{nf}}{k_f} \right) \theta'' + \frac{Pr}{2} \left(1 - \phi + \phi \left(\frac{\rho C_p_s}{\rho C_p_f} \right) \right) f \theta' = 0 \tag{9}$$

subjected to the boundary conditions of Eq. 4 which become,

$$f(0) = 0, \quad f'(0) = \lambda, \quad \theta'(0) = -Bi_x (1 - \theta(0)), \text{ at } \eta = 0 \\ f'(\eta) \rightarrow 1, \quad \theta(\eta) \rightarrow 0, \quad \text{as } \eta \rightarrow \infty \tag{10}$$

where, primes denote differentiation with respect to η , M is the magnetic parameters where $M = \sigma B_0^2 / \rho U_\infty$, Gr_x is the local Grashof number, $Gr_x = g \beta_f (T_f - T_\infty) x / U_\infty^2$ and Pr is the Prantl number, $Pr = (\mu C_p)_f / k_f$. Note that λ denotes the direction of motion of the plate. Here $\lambda = 0$ for the static plate, while $\lambda < 0$ is for the plate and the fluid moving in the opposite direction and $\lambda > 0$ indicates the plate and the fluid move in the same. To have a true similarity solution, the parameter Gr_x and Bi_x must be constant and independent of x . This condition will be satisfied if the thermal expansion coefficient β_f proportional to x^{-1} . Hence, assume of [Makinde \(2011\)](#), $\beta_f = ax^{-1}$ where a is a constant but has the appropriate dimension. Substituting $\beta_f = ax^{-1}$ into the parameter Gr_x will result $Gr = ag(T_f - T_\infty) / U_\infty^2$. For Bi_x , $h_f = bx^{-1/2}$ will gives $Bi_x = \frac{h_f}{k_{nf}} \sqrt{\frac{\nu_f}{U_\infty}}$. The quantities of engineering interest are the local skin-friction coefficient, C_f at the surface of the plate and the local Nusselt number, Nu_x which are defined as:

$$C_f = \frac{\tau_w}{\rho_f U_\infty^2}, \quad Nu_x = \frac{x q_w}{k_f (T_f - T_\infty)} \tag{11}$$

where, τ_w is the wall skin friction or shear stress at the plate and q_w is the heat flux from the plate, which is given by:

$$\tau_w = \mu_{nf} \left(\frac{\partial u}{\partial y} \right)_{y=0}, \quad q_w = -k_{nf} \left(\frac{\partial T}{\partial y} \right)_{y=0} \tag{12}$$

Substituting Eqs. 5-7 into Eq. 12 and using Eq. 11,

$$C_f(Re_x)^{\frac{1}{2}} = \frac{1}{(1-\phi)^{2.5}} f''(0), \quad \frac{Nu_x}{(Re_x)^{\frac{1}{2}}} = -\frac{k_{nf}}{k_f} \theta'(0) \quad (13)$$

3. Numerical techniques

The nonlinear ordinary differential equations in Eqs. 8–9 subject to boundary conditions of Eq. 10 are solved using the Keller–Box method, a well-known explicit finite difference scheme. This method is described in detail by [Cebeci and Bradshaw \(2012\)](#) and is very suitable for solving large numbers of nonlinear coupled equations. The following four steps are involved with this method.

1. Reduce Eqs. 8-9 to the first order system.
2. Write the difference equations using central differences.

3. Linearize the obtained algebraic equations by Newton’s method and arrange them in a matrix-vector form.
4. Solve the linear system of equations by using the block tridiagonal elimination method.

4. Results and discussion

In this section, the simulation results of different variations of flow parameters are presented and discussed. Thermophysical properties of water, kerosene and Fe₃O₄ is presented in [Table 1](#). The numerical results of non-dimensional velocity and temperatures are shown in [Figs. 2–7](#). The parameters used for simulation are $\alpha = 90^\circ$, $M = 1$, $\gamma = 45^\circ$, $Gr_x = 0.1$, $\phi = 0.05$ and $Bi_x = 0.1$, unless otherwise stated:

Table 1: Thermophysical properties of base fluids and nanoparticles ([Mojumder et al., 2015](#); [Sheikholeslami et al., 2015](#); [Ilias et al., 2017](#))

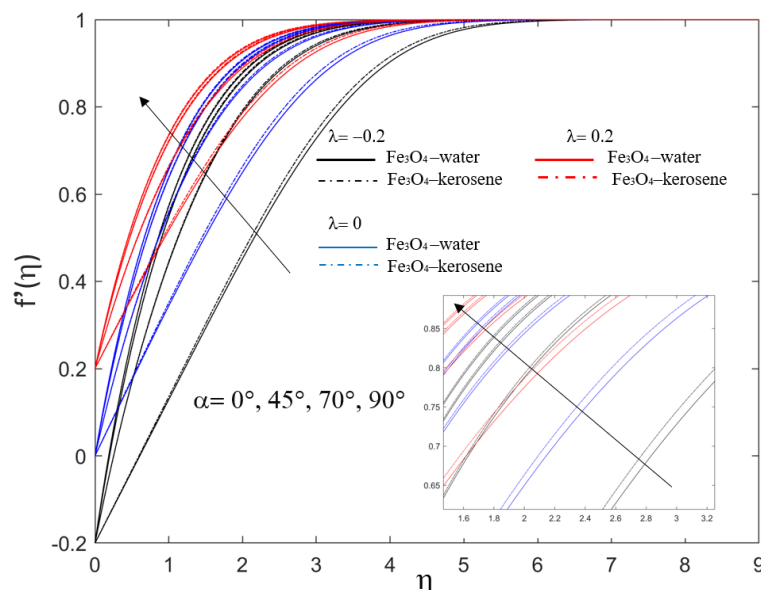
Physical Properties	Water	Kerosene	Fe ₃ O ₄
$\rho(kg/m^3)$	997.1	780	5200
$C_p(J/kgK)$	4179	2090	670
$k(W/mK)$	0.613	0.149	6
$\beta \times 10^{-5}(K^{-1})$	21	99	1.3
Pr	6.2	21	

To validate the accuracy of the numerical method, a direct comparison was made with the previously reported numerical results by [Ramesh et al. \(2016\)](#)

for different values of the Biot number. Based on [Table 2](#), the obtained results are observed to be in good agreement with the earlier findings.

Table 2: Comparison result of $\theta(0)$ for different values of the Biot number (Bi_x) when $M = 0$, $Pr = 0.72$, $Gr_x = 0.5$, $\phi = 0$ and $\frac{k_{nf}}{k_f} = 1$

Bi_x	$\alpha = 90^\circ$		$\alpha = 30^\circ$		$\alpha = 0^\circ$	
	Ramesh et al. (2016)	Present	Ramesh et al. (2016)	Present	Ramesh et al. (2016)	Present
0.05	0.1446	0.144660	0.1394	0.139477	0.1388	0.138810
0.1	0.2527	0.252756	0.2401	0.240119	0.2386	0.238622
0.2	0.4035	0.403520	0.3800	0.380015	0.3774	0.377434
0.4	0.5750	0.575012	0.5431	0.543185	0.5398	0.539854
0.6	0.6699	0.669914	0.6371	0.637147	0.6337	0.633763
0.8	0.7301	0.730168	0.6986	0.698690	0.6954	0.695454
1.0	0.7718	0.771821	0.7422	0.742245	0.7392	0.739209
5	0.9441	0.944173	0.9334	0.933433	0.9323	0.932320
10	0.9712	0.971285	0.9654	0.965434	0.9648	0.964825



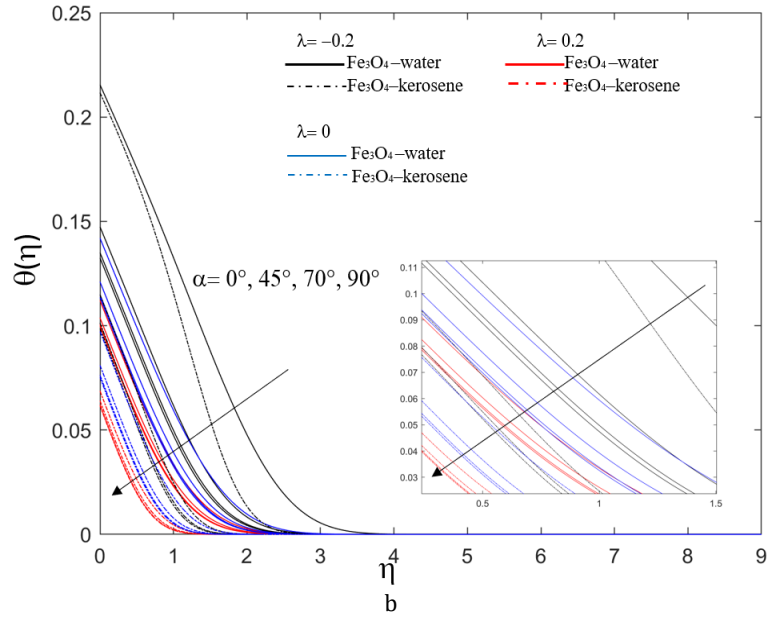


Fig. 2: Effect of aligned magnetic field angle parameter on the (a) velocity and (b) temperature profiles

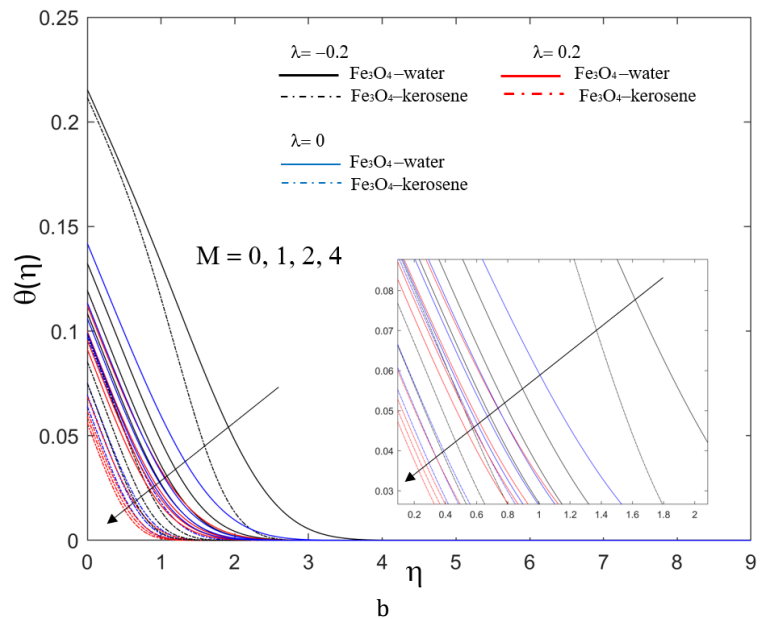
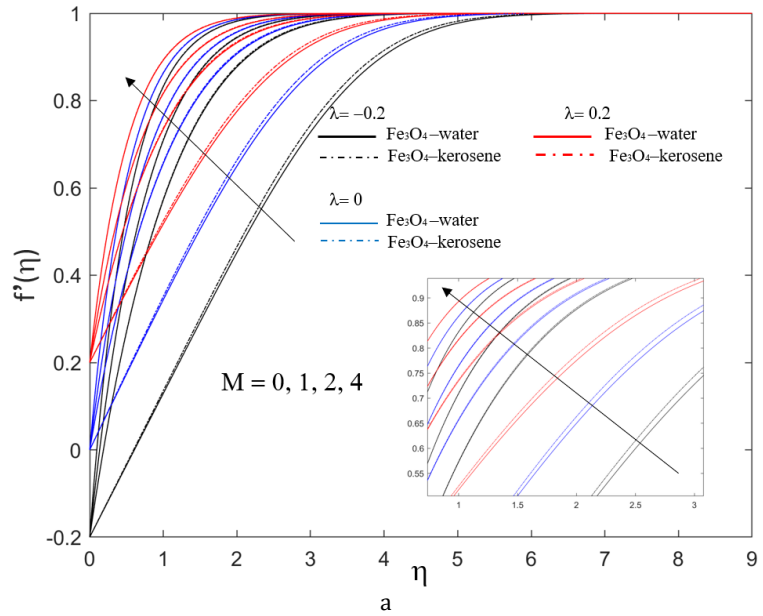


Fig. 3: Effect of magnetic strength parameter on the (a) velocity and (b) temperature profiles

As illustrated in Fig. 2, the influence of aligned magnetic field angle α , on velocity and temperature profiles of both Fe_3O_4 -water and Fe_3O_4 -kerosene magnetic nanofluids are observed. It is concluded that enhanced velocity profiles and reduced temperature profiles are achieved for both magnetic nanofluids with an increase in α . Further, a decline in the momentum boundary layer and the thermal boundary layer are seen with increasing α in the case of both magnetic nanofluids. This could be due to the strengthening of the applied magnetic field causing an increase in the value of the aligned angle ($0^\circ \leq \alpha \leq 90^\circ$) in the inclined plate. At $\alpha = 90^\circ$, this aligned magnetic field behaves like a transverse

magnetic field and the magnetic field attracts nanoparticles because of the change in the aligned magnetic field's positions. For case $\lambda = 0.2$ the velocity profiles are high but the temperature is low compared to the other two cases ($\lambda = -0.2$ and $\lambda = 0$). As observed in Fig. 3, an increase in magnetic field parameter causes a monotonic rise in the fluid velocity profiles along with a decrease in temperature. Moreover, for both magnetic nanofluids, a decrease in momentum and thermal boundary layer can be seen. From Figs. 3a and 3b, Fe_3O_4 -kerosene shows the highest velocity profiles and lowest temperature compared to Fe_3O_4 -water.

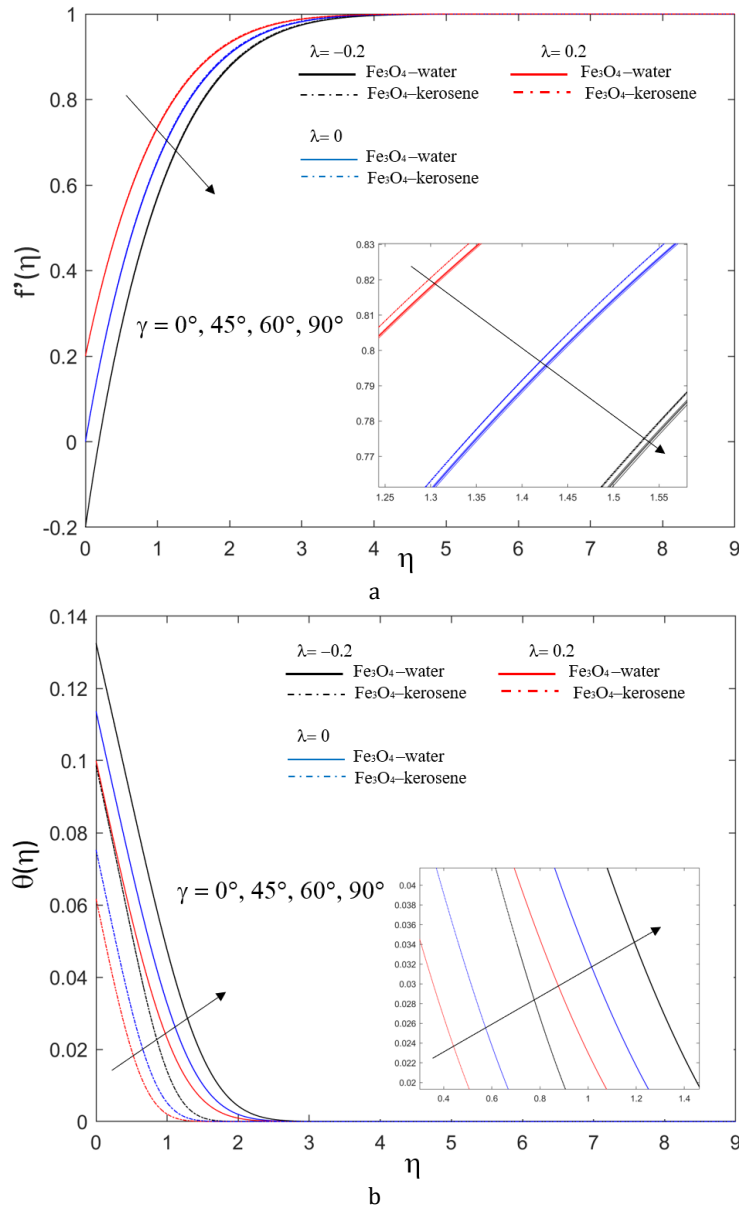


Fig. 4: Effect of inclined plate angle parameter on the (a) velocity and (b) temperature profiles

Fig. 4 shows a deceleration occurring in the velocity of magnetic nanofluids with an increase in the inclination of plate γ . The plate assumes a vertical position if $\gamma = 0^\circ$ while it is horizontal if $\gamma = 90^\circ$. For $\gamma = 90^\circ$, the gravitational effect is minimum while it is maximum for $\gamma = 0^\circ$. An increase of γ causes increase in the momentum and thermal

boundary layer for both magnetic nanofluids. Fig. 5, shows that an increase in the Gr_x , the velocity are increases while the temperature profile decrease. As can be seen from Figs. 5a and 5b, the thicknesses of thermal and momentum boundary layers decrease with a rise in Gr_x due to the buoyancy effect. Positive values of local Grashof number $Gr_x > 0$ are

employed for the computations, which corresponds to the cooling issue with regards to the application.

Most engineering applications often encounter cooling issue.

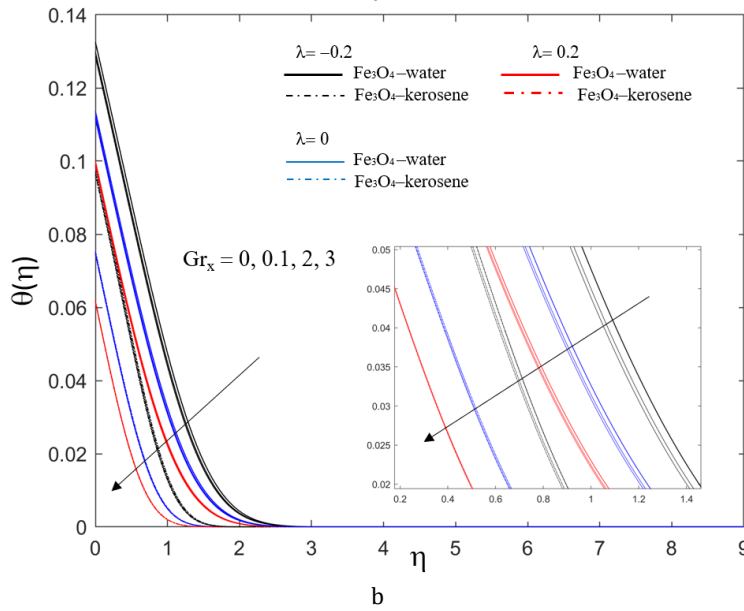
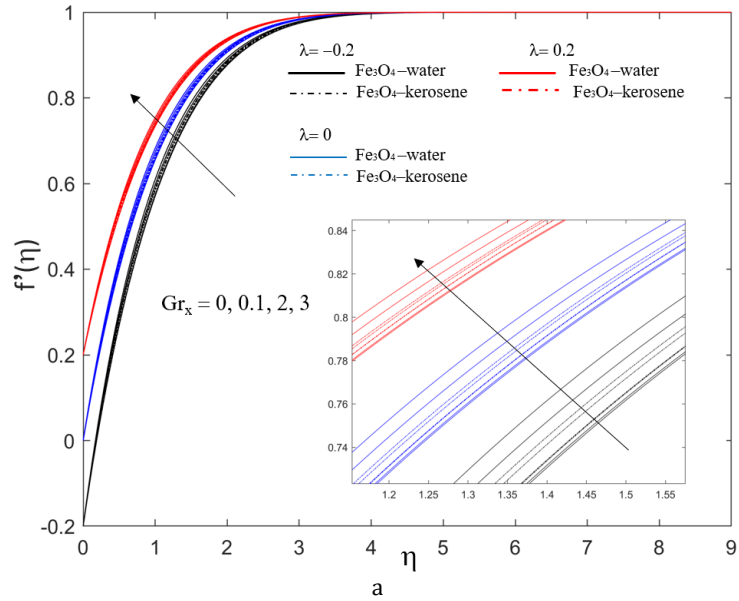
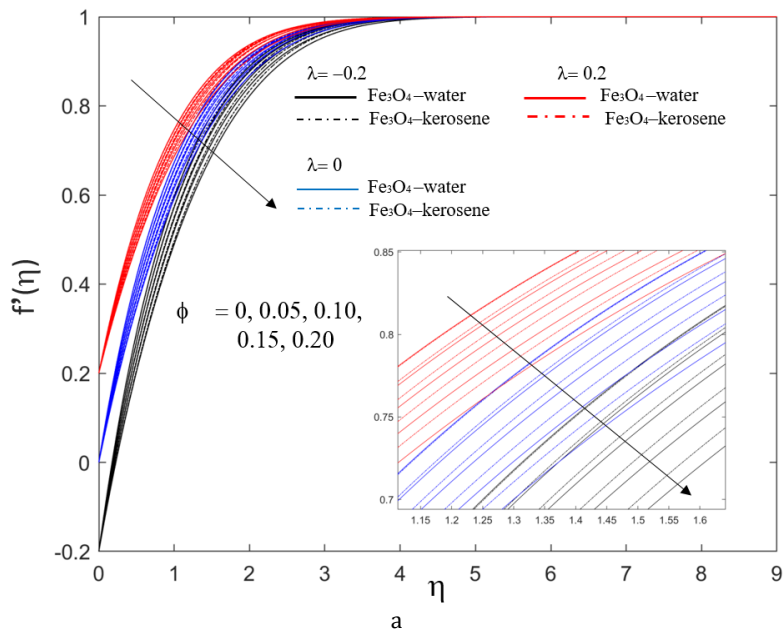


Fig. 5: Effect of local Grashof number on the (a) velocity and (b) temperature profiles



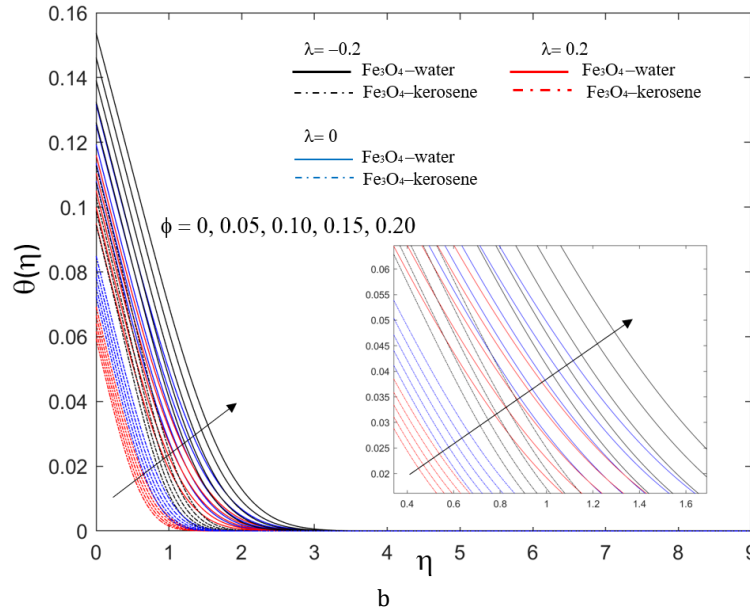


Fig. 6: Effect of volume fraction of nanoparticles on the (a) velocity and (b) temperature profiles

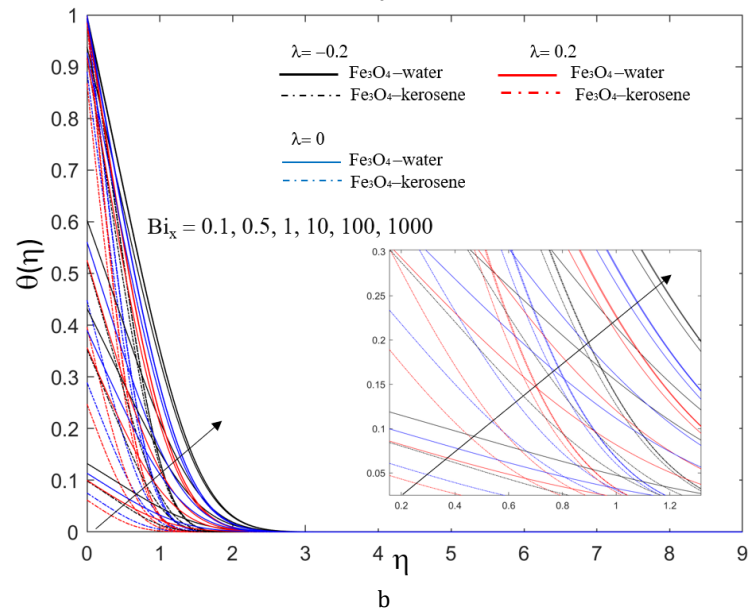
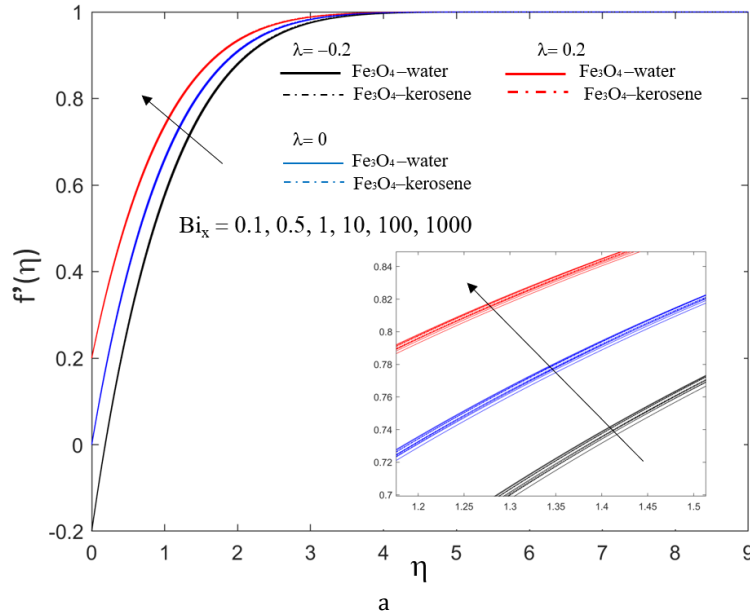


Fig. 7: Effect of convection parameter, Bi_x on the (a) velocity and (b) temperature profiles for different values of the plate velocity parameter (λ)

Fig. 6 presents the velocity and temperature profiles for different volume fraction values, ϕ . For both magnetic nanofluids, the velocity decreases with an increase in ϕ , while the momentum boundary layer increases. This is caused by the increase in the number of collisions amongst solid particles, which consequently results in the velocity reduction of magnetic nanofluids. For temperature distribution, temperature, as well as the thermal boundary layer thickness, are enhanced with an increase in solid ϕ . For the temperature profile, as the distance increases from the boundary, the distribution asymptotically moves toward zero. This is in agreement with the physical behavior that states an increase in thermal conductivity as well as thickness of the thermal boundary layer with the increase in the volume fraction of Fe_3O_4 -kerosene and Fe_3O_4 -water.

Fig. 7 illustrates the effect of the convection parameter, Bi_x on velocity and temperature profiles for both magnetic nanofluids at different values of moving inclined plate parameters, $\lambda = -0.2$, $\lambda = 0$, and $\lambda = 0.2$. It is seen from Figs. 7a and 7b that magnetic nanofluids' velocity increases with the increase of Bi_x , while the momentum boundary layer decreases. It can also be concluded that the temperature profiles and thermal boundary layer thickness increase with the increase of Bi_x . A similar observation has been reported by many reserachers (Ramesh et al., 2016; Sharma and Gupta, 2017; Dash and Mishra, 2022; Yesodha et al., 2022). As the value of the parameter Bi_x increases, the intensity of convective heating on the plate surface increases, which leads to an increasing rate of convective heat transfer from the hot fluid on the lower surface of the plate to the magnetic nanofluids on the upper surface. The graph also reveals that temperature

increases rapidly near the surface due to the increasing values of Bi_x . Thermal boundary layer for Fe_3O_4 -water is higher compared to Fe_3O_4 -kerosene. For $\lambda = 0.2$, the velocity and temperatue profiles is higher compared to another two cases. For each curve, the vertical intercept gives the plate surface temperature. The plate surface temperature increases as Bi_x increases. As $Bi_x \rightarrow \infty$, the solution approaches the classical solution for the constant surface temperature. This can be seen from the boundary condition 10 which reduces to $\theta'(0) - Bi_x(1 - \theta) \rightarrow \infty$ as $Bi_x \rightarrow \infty$.

Table 3 presents the variation occurring in skin friction coefficient for different $\alpha, M, \gamma, Gr_x, \phi$ and Bi_x values applicable to various inclined plate movements. For all three cases, it could be clearly seen that parameters ($\alpha, M, \gamma, Gr_x, \phi$ and Bi_x) had a boosting effect on the friction factor coefficients, especially for the case when $\lambda = -0.2$. The highest value could be seen for the case when M, Gr_x and ϕ are modified. It is interesting to note that in the case $\lambda = 0.2$, the friction factor coefficient was found to be lesser than in the cases of $\lambda = -0.2$ and $\lambda = 0$. So, the friction between the particles can be decreased in the case of $\lambda = 0.2$. In all cases, skin friction for Fe_3O_4 -kerosene is higher than Fe_3O_4 -water except for γ .

Table 4 lists the different variations in the Nusselt number for different values of dimensionless parameters. For all three cases, an enhancement was seen in the heat transfer rates with improving values of $\alpha, M, \gamma, Gr_x, \phi$ and Bi_x , especially for $\lambda = 0.2$ can be benefit to the heat transfer rates. The highest value could be seen for the case when ϕ and Bi_x in the case $\lambda = 0.2$. In all cases, Nusselt number for Fe_3O_4 -kerosene is higher than Fe_3O_4 -water except for γ .

Table 3: Variation in local skin friction for both magnetic nanofluids at different dimensionless parameters

α	M	γ	Gr_x	ϕ	Bi_x	Magnetic nanofluids (Fe_3O_4 -water)			Magnetic nanofluids (Fe_3O_4 -kerosene)		
						$\lambda = -0.2$	$\lambda = 0$	$\lambda = 0.2$	$\lambda = -0.2$	$\lambda = 0$	$\lambda = 0.2$
0°						0.376803	0.397854	0.372688	0.383547	0.404891	0.380414
45°	1	45°	0.1	0.05	0.1	0.972052	0.841168	0.697534	0.973292	0.843750	0.700857
70°						1.252000	1.066931	0.872440	1.252722	1.068702	0.874792
90°						1.325944	1.127130	0.919493	1.326577	1.128750	0.921657
						0.376803	0.397854	0.372688	0.383547	0.404891	0.380414
90°	1	45°	0.1	0.05	0.1	1.325944	1.127130	0.919493	1.326577	1.128750	0.921657
	2					1.841523	1.550241	1.252808	1.841800	1.551207	1.254140
	4					2.581253	2.162081	1.738587	2.581357	2.162658	1.739403
						1.327898	1.128664	0.920739	1.327710	1.129485	0.922183
90°	1	45°	0.1	0.05	0.1	1.325944	1.127130	0.919493	1.326577	1.128750	0.921657
		60°				1.324560	1.126044	0.918611	1.325775	1.128230	0.921286
		90°				1.321214	1.123420	0.916481	1.323837	1.126973	0.920388
						1.321214	1.123420	0.916481	1.323837	1.126973	0.920388
90°	1	45°	0	0.05	0.1	1.325944	1.127130	0.919493	1.326577	1.128750	0.921657
			2			1.412744	1.196119	0.975921	1.377303	1.162133	0.945653
			3			1.456341	1.231363	1.005041	1.403068	1.179429	0.958183
						1.236444	1.047452	0.851737	1.234670	1.045716	0.850166
90°	1	45°	0	0.1	0.1	1.325944	1.127130	0.919493	1.326577	1.128750	0.921657
			0.1			1.426870	1.217024	0.995945	1.430301	1.222515	1.002384
			0.15			1.541411	1.319088	1.082754	1.548088	1.329037	1.094081
			0.2			1.672347	1.435805	1.182029	1.682795	1.450903	1.198965
90°	1	45°	0.1	0.05	0.1	1.325944	1.127130	0.919493	1.326577	1.128750	0.921657
					0.5	1.336623	1.136148	0.927223	1.333635	1.133792	0.925477
					1	1.342699	1.141712	0.932303	1.338297	1.137543	0.928544
					10	1.354573	1.153616	0.944052	1.349183	1.147937	0.938226
90°	1	45°	0.1	0.05	100	1.356531	1.155723	0.946268	1.351255	1.150229	0.940632
					1000	1.356740	1.155950	0.946510	1.351481	1.150486	0.940909

Table 4: Variation in local Nusselt number for both magnetic nanofluids at different dimensionless parameters

α	M	γ	Gr_x	ϕ	Bi_x	Magnetic nanofluids (Fe ₃ O ₄ -water)			Magnetic nanofluids (Fe ₃ O ₄ -kerosene)			
						$\lambda = -0.2$	$\lambda = 0$	$\lambda = 0.2$	$\lambda = -0.2$	$\lambda = 0$	$\lambda = 0.2$	
0°						0.087569	0.095772	0.099109	0.090360	0.103473	0.106736	
45°	1	45°	0.1	0.05	0.1	0.095167	0.098106	0.100032	0.101451	0.105329	0.107278	
70°						0.096549	0.098789	0.100380	0.103012	0.105870	0.107495	
90°						0.096834	0.098941	0.100462	0.103320	0.105991	0.107546	
						0.087569	0.095772	0.099109	0.090360	0.103473	0.106736	
	1	45°	0.1	0.05	0.1	0.096834	0.098941	0.100462	0.103320	0.105991	0.107546	
90°	2					0.098285	0.099774	0.100938	0.104827	0.106652	0.107855	
	4					0.099500	0.100547	0.101419	0.106022	0.107272	0.108182	
						0.096841	0.098944	0.100464	0.103324	0.105992	0.107547	
	1	0°	0.1	0.05	0.1	0.096834	0.098941	0.100462	0.103320	0.105991	0.107546	
90°		45°				0.096829	0.098939	0.100460	0.103318	0.105990	0.107546	
		60°				0.096817	0.098933	0.100457	0.103311	0.105988	0.107545	
		90°				0.096817	0.098933	0.100457	0.103311	0.105988	0.107545	
	1	45°	0.1	0.05	0.1	0.096834	0.098941	0.100462	0.103320	0.105991	0.107546	
90°						2	0.097128	0.099092	0.100546	0.103486	0.106042	0.107566
						3	0.097267	0.099167	0.100588	0.103567	0.106068	0.107577
							0.087391	0.089198	0.090504	0.090550	0.092764	0.094065
	1	45°	0.1	0	0.1	0.096834	0.098941	0.100462	0.103320	0.105991	0.107546	
90°				0.05		0.106895	0.109352	0.111120	0.117277	0.120495	0.122352	
				0.15		0.117634	0.120503	0.122560	0.132590	0.136473	0.138688	
				0.2		0.129118	0.132475	0.134870	0.149466	0.154163	0.156808	
	1	45°	0.1	0.1	0.1	0.096834	0.098941	0.100462	0.103320	0.105991	0.107546	
90°				0.5		0.316907	0.340425	0.359012	0.370852	0.407470	0.431387	
				1		0.443003	0.490123	0.529456	0.548681	0.632448	0.691855	
				10		0.691267	0.812167	0.925370	0.967164	1.258666	1.516406	
	1	45°	0.1	0.05	0.1	100	0.732456	0.869440	1.000313	1.047349	1.397321	
				100		0.732456	0.869440	1.000313	1.047349	1.397321	1.721888	
				1000		0.736849	0.875617	1.008483	1.056111	1.412893	1.745549	

5. Conclusions

A numerical study of the MHD free convection boundary layer flow for both water and kerosene magnetic nanofluids through an inclined plate subjected to a magnetic field has been analyzed. The following conclusions are derived.

1. The velocity increases with an increase of aligned magnetic field angle α , magnetic strength M , Grashof number Gr_x and Biot number Bi_x .
2. The velocity decreases with an increase of inclined plate angle γ and volume fractions of nanoparticles ϕ .
3. The temperature increases with an increase of inclined plate angle γ , volume fraction of nanoparticles ϕ , and Biot number Bi_x .
4. The temperature decreases with an increase of aligned magnetic field angle α , magnetic strength M and Grashof number Gr_x .
5. Skin friction coefficient and Nusselt number increase with the increase of aligned magnetic field angle α , magnetic strength M , Grashof number Gr_x , the volume fraction of nanoparticles ϕ and Biot number Bi_x for all three cases of inclined plate movements, $\lambda = -0.2$, $\lambda = 0$, and $\lambda = 0.2$.
6. Skin friction coefficient and Nusselt number decrease with the increase of inclined plate angle γ for all three cases of inclined plate movements, $\lambda = -0.2$, $\lambda = 0$, and $\lambda = 0.2$.

List of symbols

- α Aligned angle of magnetic field
- α_{nf} Thermal diffusivity of magnetic nanofluids
- β_f Thermal expansion coefficient
- γ Plate inclination angle

- η Boundary layer thickness
- $\theta(\eta)$ Non - dimensional temperature function
- λ Velocity ratio parameter
- μ_f Dynamic viscosity of base fluid
- μ_{nf} Dynamic viscosity of magnetic nanofluids
- ρ Density
- ρ_f Density of base fluid
- ρ_s Density of nanoparticles
- ρ_{nf} Density of magnetic nanofluids
- $(\rho C_p)_{nf}$ Heat capacity of magnetic nanofluids
- $(\rho C_p)_f$ Heat parameters of base fluid
- $(\rho C_p)_s$ Heat parameters of nanoparticles
- $(\rho\beta)_{nf}$ Thermal expansion of magnetic nanofluids coefficient
- σ Electrical conductivity
- τ_w Wall shear stress
- ϕ Nanoparticles volume fraction
- $\psi(x, y)$ Stream function
- $B(x)$ Transverse magnetic field
- B_0 Magnetic field strength
- Bi_x Biot number
- $f(\eta)$ Non - dimensional stream function
- C_f Local skin - friction coefficient
- Gr_x Local Grashof number
- g Gravitational acceleration
- h_f Heat transfer coefficient
- k_f Thermal conductivity of base fluid
- k_s Thermal conductivity of nanoparticles.
- k_{nf} Thermal conductivity of magnetic nanofluids
- M Magnetic strength parameter
- Nu_x Local Nusselt number
- Pr Prantl number
- q_w Heat flux
- Re_x Reynolds number
- T Temperature
- T_f Temperature of hot fluid
- T_∞ Temperature at the free stream
- $U_w(x)$ Plate velocity
- U_∞ Velocity in free stream
- u Velocity in x -direction

v	Velocity in y -direction
x	Dimensionless coordinate axis along the inclined plate
y	Dimensionless coordinate axis normal to the surface plate

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Compliance with ethical standards

Conflict of interest

The author(s) declared no potential conflicts of interest with respect to the research, authorship, and/or publication of this article.

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